

# On the statistical stability of families of attracting sets and the contracting Lorenz attractor

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Dynamical Systems from a Pacific(o) point of view,  
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## Ergodic theory in a (tiny) nutshell

**Dynamics:**  $f : M \rightarrow M$ **Observable:**  $\varphi : M \rightarrow \mathbb{R}$ 

- **Invariant measure:**  $\mu(f^{-1}A) = \mu(A)$ ;
- **Ergodic measure:**  $A = f^{-1}A \implies \mu(A) \in \{0, 1\}$ .
- **Birkhoff Ergodic Theorem:** if  $\mu$  is ergodic, then

$$\underbrace{\text{Time average}} \quad \text{tends to} \quad \underbrace{\text{space average}}$$

$$\varphi_n(x) = \frac{1}{n} \sum_{i=0}^{n-1} \varphi(f^i x) \quad \xrightarrow[n \rightarrow \infty]{\mu\text{-a.e.}} \quad \mathbb{E}(\varphi) = \int \varphi d\mu.$$

# Continuous time

$M$  smooth Riemannian manifold

$X^t : M \rightarrow M$  smooth flow (i.e.,  $X^{t+s} = X^t \circ X^s$  for  $s, t \in \mathbb{R}$ )

- **Invariant measure:**  $\mu(X^t A) = \mu(A)$ ,  $\forall 0 < t \leq 1$ ;
- **Ergodic measure:**  
 $\exists \varepsilon > 0 : A = X^t A, \forall 0 < t < \varepsilon \implies \mu(A) \in \{0, 1\}$ .
- **Birkhoff Ergodic Theorem:** if  $\mu$  is ergodic, then

$$\underbrace{\varphi_T(y) = \frac{1}{T} \int_0^T \varphi(X^t y) dt}_{\text{Time average}} \xrightarrow[T \rightarrow \infty]{\mu\text{-a.e.}} \underbrace{\mathbb{E}(\varphi) = \int \varphi d\mu}_{\text{space average}}$$

# Physical/SRB measure.

Is there an **invariant physical/SRB measure**  $\mu_{SRB}$ ?

That is, a measure  $\mu_{SRB}$  so that  $Leb(B(\mu_{SRB})) > 0$  where

$$B(\mu_{SRB}) = \left\{ y \in M : \frac{1}{T} \int_0^T \varphi(X^t y) dt \xrightarrow{T \rightarrow \infty} \int \varphi d\mu_{SRB}, \forall \varphi \in C(M) \right\}$$

is the **ergodic basin** of  $\mu_{SRB}$ .

**This kind of measure provides asymptotic information on a set of trajectories that one hopes is large enough to be observable in real-world models.**

# Statistical stability

Can we **allow for small errors on the formulation of the dynamics not to disturb too much** the long term behavior?

If we consider

$\mathcal{U} = \{Y \text{ vector field s.t. } Y^t \text{ admits physical measure } \mu^Y\}$

and  $X \in \mathcal{U}$ . Then  $X$  is **statistically stable** if

$$\underbrace{\|Y_n - X\|_{C^1}}_{\text{Small errors}} \xrightarrow[n \rightarrow \infty]{Y_n \in \mathcal{U}} 0 \quad \text{disturb} \quad \implies \quad \int \varphi d\mu^{Y_n} \xrightarrow[n \rightarrow \infty]{\varphi \in C^0(M)} \int \varphi d\mu \quad \underbrace{\text{space average}}$$

# Some known results on existence of physical measures

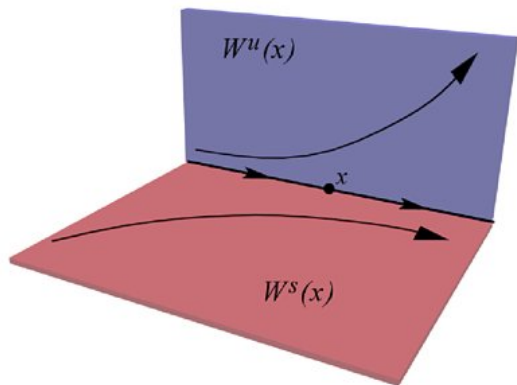
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## Hyperbolic versus singular flows

# Hyperbolic flows

**Hyperbolic flows:** all trajectories have a pair of complementary directions:

- in one of them all orbits converge to the trajectory;
- in the other direction all orbits diverge from the trajectory.



# Hyperbolic flows are “classical”

**Hyperbolic Theory is the basis for Dynamical Systems Theory:** it provides the most mathematically rigorous and deep understanding of an important class of dynamical systems.

**This is an *open class of flows*:** all flows nearby an hyperbolic flow are also hyperbolic.

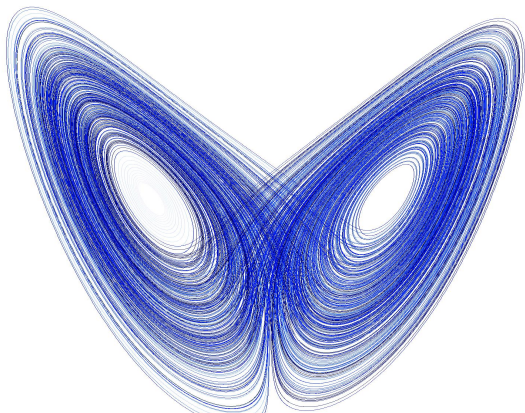
**Hyperbolic flows do not admit fixed points (singularities or equilibria) accumulated by regular orbits in invariant sets.**

However, **there are important open classes of systems which are not hyperbolic** and that frequently appear in applications.



# Singular flows which are “almost hyperbolic”

The Lorenz attractor is a flow with an equilibrium accumulated by regular orbits which also belongs to an open class



# Attractors and attracting sets

An invariant compact set  $\Lambda$  is an **attracting set** for a vector field  $X$  if there exists a **trapping region**  $U$  s.t.

$$\overline{X^t(U)} \subset U \text{ for large } t > 0 \text{ and } \Lambda = \bigcap_{t>0} \overline{X^t(U)}.$$

An attracting set becomes an **attractor** if  $\Lambda$  is transitive, that is, we can find  $x \in \Lambda$  s.t.

$$\mathcal{O}^+(x) = \{X^t(x) : t > 0\} \text{ is dense in } \Lambda.$$

## Existence of physical measures

family vec. fields	physical measures	ergodic basins
Anosov flows (transitive) Axiom A flows (Hyperbolic)	unique one for each attractor	$\text{Leb}(M \setminus B(\mu)) = 0$ $\text{Leb}(U_i \setminus B(\mu_i)) = 0$ for each attractor
geometric Lorenz attractor	unique	$\text{Leb}(U \setminus B(\mu)) = 0$
contracting Lorenz attractor	unique	$\text{Leb}(U \setminus B(\mu)) = 0$
sectional- hyperbolic attracting sets	finitely many	$\text{Leb}(U \setminus \cup_i B(\mu_i)) = 0$

**Except the contracting Lorenz (Rovella) attractor, all the other families are  $C^r$  open families ( $r \geq 1$ ).**

# Physical measures and equilibrium states

# Physical measures and equilibrium states

Let  $\Lambda$  be a **sectional-hyperbolic attracting set / geometrical or contracting Lorenz attractor / hyperbolic attractor** for a  $C^2$  vec. field  $G$  with trapping region  $U$ . Then the following are equivalent:

- 1  $\mu$  is an equilibrium state with respect to the central jacobian:  $h_\mu(X^1) = \int \log |\det DX^1|_{E^c}| d\mu > 0$ ;
- 2  $\mu$  is a *SRB* measure, i.e., admits an **abs. cont. disintegration along unstable manifolds**;
- 3  $\mu$  is a physical measure, i.e.,  $\text{Leb}(B(\mu)) > 0$ ;

Moreover, the family  $\mathbb{E}$  of all invariant physical measures is the following convex hull

$$\mathbb{E} = \left\{ \sum_{i=1}^k t_i \mu_i : \sum_i t_i = 1; 0 \leq t_i \leq 1, i = 1, \dots, k \right\}.$$

# Entropy expansiveness

# Topological entropy

Let  $g : M \rightarrow M$  be a continuous map and  $K \subset M$ .  
For  $\varepsilon > 0$ ,  $n \geq 1$  and  $x \in M$

$$B(x, \varepsilon, n) = \{y \in M : d(g^j x, g^j y) < \varepsilon, \quad \forall 0 \leq j < n\}.$$

A subset  $E \subset M$  is a  $(n, \varepsilon)$ -**generator for**  $K$  if  $\{B(y, \varepsilon, n) : y \in E\}$  is an open cover of  $K$ .

Let  $r_n(K, \varepsilon)$  be the smallest cardinality of a  $(n, \varepsilon)$ -generator for  $K$  and  $r(K, \varepsilon) = \limsup \log r_n(K, \varepsilon)^{1/n}$ .

The **topological entropy of  $g$  on  $K$**  is given by

$$h_{\text{top}}(g, K) = \lim_{\varepsilon \rightarrow 0} r(K, \varepsilon),$$

and the **topological entropy of  $g$**  is defined by  $h_{\text{top}}(g) = h_{\text{top}}(g, M)$ .

# Entropy expansiveness

For  $x \in M$  and  $\varepsilon > 0$  we define the **two-sided  $\varepsilon$ -dynamical ball at  $x$**  as

$$B(x, \varepsilon, \infty) = \{y \in M : d(g^n x, g^n y) < \varepsilon, \forall n \in \mathbb{Z}\}$$

and say that  $g$  **is  $\varepsilon$ -entropy expansive** if all these infinite dynamical balls have zero topological entropy, that is,

$$\sup_{x \in M} h_{top}(g, B(x, \varepsilon, \infty)) = 0.$$

## Proposition

If  $G$  is entropy expansive, then the metric entropy function  $\mu \mapsto h_\mu(X^1)$  is **upper semicontinuous**.



# Conditions for statistical stability — — — of families of vector fields

# Conditions for statistical stability

Let  $\mathcal{G}$  be a collection of vec. fields with a trapping region  $U$ ,  $s \in N \mapsto G_s \in \mathcal{G}$  is a continuous parametrization, and  $\Lambda_s(U) = \bigcap_{t>0} \overline{\phi_t^{G_s}(U)}$  the corresponding attracting set.

## Theorem

Assume each  $\Lambda_s$  supports finitely many ergodic physical measures  $\mu_i^s$ ,  $1 \leq i \leq k_s$  s.t.  $\text{Leb}(U \setminus \sum_i B(\mu_i^s)) = 0$ ; and

- 1 there are potentials  $\psi_s : \Lambda_s \rightarrow \mathbb{R}$  s.t.  
 $0 = h_\mu(\phi_1^{G_s}) + \int \psi_s d\mu \iff \mu$  a physical measure;
- 2 the following map is continuous  
 $\Psi : W(U) := \{(s, x) \in N \times U : x \in \Lambda_s(U)\} \rightarrow \mathbb{R}$  given by  
 $\Psi(s, x) = \psi_s(x)$ ; and
- 3 the family  $\mathcal{G}$  is robustly entropy expansive.

Then  $G_s \in \mathcal{G}$  is statistically stable.

# Statistical stability with several physical measures

In the general setting of the statement we have:

## Statistical stability

For each converging sequence  $s_n \in N$  to  $s \in N$  and every choice  $\mu^{s_n}$  of a physical measure supported on  $\Lambda_{s_n}$ , **every weak\* accumulation point  $\mu$  of  $(\mu^{s_n})_{n \geq 1}$  is a convex linear combination of the ergodic physical measures of  $\Lambda_s$ .**

More precisely, this means that

- 1 there are weights  $\alpha_i \geq 0$  such that  $\sum_i \alpha_i = 1$  and  $\mu = \sum_i \alpha_i \mu_i^s$ ; and
- 2 we have  $\left| \int \varphi d\mu^{s_n} - \sum_i \alpha_i \int \varphi d\mu_i^s \right| \xrightarrow{n \rightarrow \infty} 0$  for every continuous observable  $\varphi : U \rightarrow \mathbb{R}$ .

# Known examples of application

# Axiom A (hyperbolic) flows

Attracting sets for an Axiom A vector field  $G$  of class  $C^2$  are finite unions of hyperbolic attractors (basic pieces) which admit a unique physical measure that is also the unique eq. state w.r.t.  $\psi_G = -\log |\det D\phi_1^G| |E^u|$  (**here  $\phi_t^G$  is the flow generated by  $G$** ).

**Moreover, the hyperbolic property is robust and so the map  $X \mapsto \psi_X$  is well defined for vector fields  $X$  close to  $G$  in the  $C^2$  topology.**

In addition, each  $X$  is a neighborhood of  $G$  is not only hyperbolic but also entropy expansive.

**Hence, we may apply the Main Theorem and deduce statistical stability for each basic piece of an Axiom A vector field which is an attractor.**

# Singular-hyperbolic attracting sets – encompassing the geometrical Lorenz attractor

Singular-hyperbolicity is an extension of the notion of hyperbolicity encompassing sets with equilibria accumulated by regular orbits.

A **singular-hyperbolic set**  $\Lambda$  is

- a *partially hyperbolic set*: there exists a splitting  $T_\Lambda M = E^s \oplus E^{cu}$ , where  $d_s = \dim E_x^s \geq 1$  and  $d_{cu} = \dim E_x^{cu} = 2$  for  $x \in \Lambda$ , and constants  $C > 0$ ,  $\lambda \in (0, 1)$  s.t. for  $t > 0$  we have
  - *uniform contraction on  $E^s$* :  $\|D\phi_t|E_x^s\| \leq C\lambda^t$ ; and
  - *domination*:  $\|D\phi_t|E_x^s\| \cdot \|D\phi_{-t}|E_{\phi_t x}^{cu}\| \leq C\lambda^t$ .
- with *area expansion on  $E^{cu}$* :  $|\det(D\phi_t|E_x^{cu})| \geq C\lambda^{-t}$ ;
- any equilibrium of  $\Lambda$ , if any, is hyperbolic.

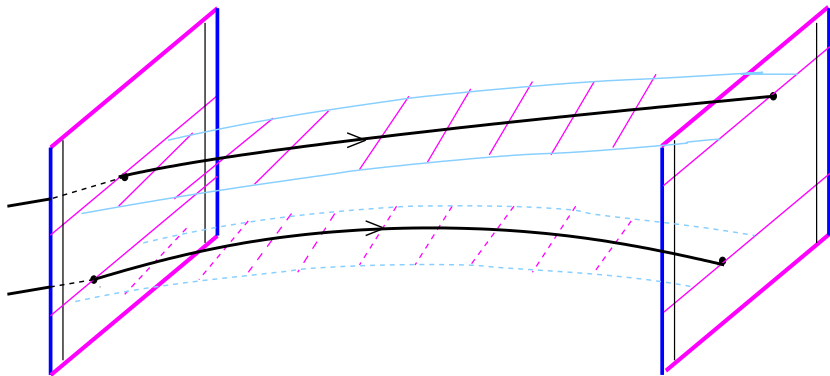
# Singular-hyperbolic attracting sets and statistical stability

The assumptions of the Main Theorem are known to hold for singular-hyperbolic attracting sets with the potential  $\psi_G = -\log |\det D\phi_1^G|_{E^{cu}}$ : **existence of finitely many physical/SRB measures** and **robust entropy expansiveness** for singular-hyperbolic attracting sets is established in

- A., **M. J. Pacifico**, Pujals, Viana: “Singular-hyperbolic attractors are chaotic”. TAMS, 2009. [unique SRB, transitive case]
- A., Souza, Trindade: “Upper Large Deviations Bound for Singular-Hyperbolic Attracting Sets”. JDDE, 2019. [finite # erg. SRB, non-transitive]
- **M. J. Pacifico**, F. Yang, J. Yang: “Entropy theory for sectional hyperbolic flows”. An. I’IHP, 2020. [robust entropy expansiveness]

# Other strategies to obtain statistical stability

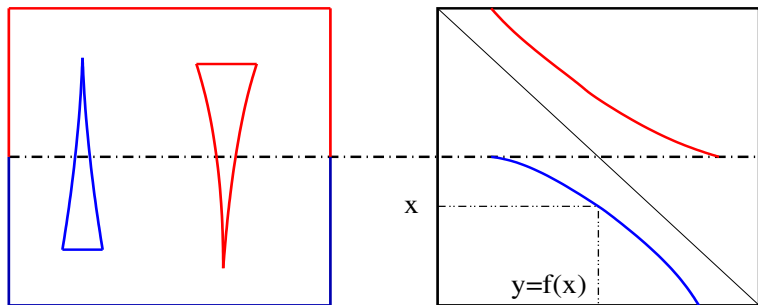
The dynamics on singular-hyperbolic attracting sets is amenable to **reduction to a global Poincaré return map on a finite collection of cross-sections**:





# Reduction to one-dimensional transformation

There is also a further reduction to **a quotient map along the stable leaves tangent to the stable bundle**.



For the geometric **Lorenz attractor**, this procedure ends with the **one-dimensional Lorenz transformation**.

# Statistical properties from the reduction

**The physical measure can be constructed from the acip of the one-dimensional map and statistical stability can be deduced** from this:

- Alves, Soufi: “Statistical stability of geometric Lorenz attractors”. Fund. Math., 2014.
- Bahsoun, Ruziboev: “On the statistical stability of Lorenz attractors with a  $C^{1+\alpha}$  stable foliation”. ETDS, 2018.

Many other finer properties can be deduced:

- A., Melbourne: “Exponential decay of correlations for nonuniformly hyperbolic flows with a  $C^{1+\alpha}$  stable foliation, including the classical Lorenz attractor”. AHP, 2016.
- Bahsoun, Melbourne, Ruziboev: “Variance Continuity for Lorenz Flows”. AHP, 2020

# New examples of application

# Sectional-hyperbolic flows

Sectional-hyperbolicity is an extension of singular hyperbolicity with central dimension  $d_{cu} > 2$  where the area expansion property is replaced by **sectional expansion**: there are  $K, \theta > 0$  s.t. **for every two-dimensional subspace**  $P_x \subset E_x^{cu}$

$$|\det(D\phi_t | P_x)| \geq Ke^{\theta t} \quad \text{for all } x \in \Lambda, t \geq 0.$$

## Theorem (A., ETDS, 2021)

Every sectional-hyperbolic attracting set for a  $C^2$  vector field  $G$  admits finitely many  $\mu_1, \dots, \mu_k$  ergodic physical/SRB measures s.t.  $h_{\mu_i}(\phi_1^G) + \int \psi_G d\mu_i = 0$  and  $\text{Leb}(U \setminus \sum_j B(\mu_j)) = 0$ .

Recall that  $\psi_G = -\log |\det D\phi_1^G | E^{cu}|$ .

# Statistical stability for sectional-hyperbolic attracting sets

Since

- **sectional-hyperbolicity is a  $C^1$  open property**, then the family of vector fields having sectional-hyperbolic attracting sets is  $C^1$  open; and
- **entropy expansiveness** was already obtained by
  - **M. J. Pacifico**, F. Yang, J. Yang: “Entropy theory for sectional hyperbolic flows”. An. I’HP, 2020.

then we have **all the conditions for statistical stability for sectional-hyperbolic attractors**.

# Contracting Lorenz family of attractors

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also known as “Rovella attractors”

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which is **NOT AN OPEN FAMILY** of vector fields

# The Rovella family of attractors

This is a **modification of the geometric Lorenz attractor** – the **expanding direction at the equilibrium is replaced by an area contracting direction**:

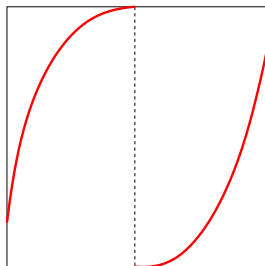
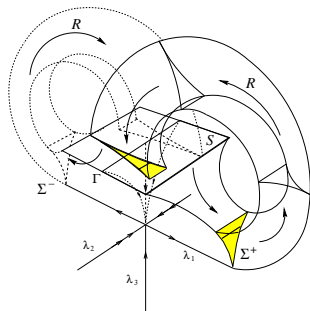
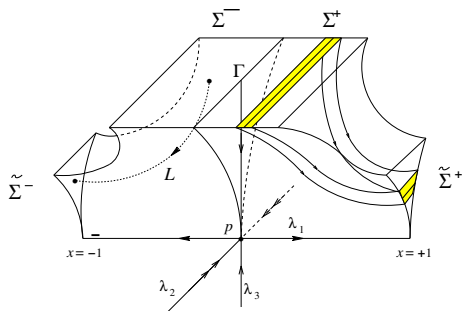
start with a linear vector field

$(\dot{x}, \dot{y}, \dot{z}) = (\lambda_1 x, \lambda_2 y, \lambda_3 z)$  in  $[-1, 1]^3$  with real eigenvalues at the singularity s.t.

$$-\lambda_2 > -\lambda_3 > \lambda_1 > 0, \quad r = -\frac{\lambda_2}{\lambda_1}, \quad s = -\frac{\lambda_3}{\lambda_1}, \quad \text{and} \quad r > s + 3.$$

**Note that  $\lambda_1 + \lambda_3 < 0$  while in the geometric Lorenz attractor the construction starts with  $\lambda_1 + \lambda_3 > 0$ .**

# Geometric construction and quotient map





# Smooth foliation and quotienting

The condition  $r > s + 3$  ensures the existence of a  $C^3$  **uniformly contracting stable foliation for the Poincaré first return map of all small enough perturbations of the contracting geometric Lorenz flow.**

Using this, we write the Poincaré first return map as  $R_0(x, y) = (T_0(x), H_0(x, y))$ , where  $H_0(x, y)$  uniformly contracts distances along  $y$  and

- 1  $T_0 : [-1/2, 1/2] \setminus \{0\} \rightarrow [-1/2, 1/2]$  is piecewise  $C^3$  with two onto branches s.t.  $T'_0(x) = O(x^{s-1})$  at  $x = 0$ ;
- 2  $T_0(0^+) = -1/2$  and  $T_0(0^-) = +1/2$ ;
- 3  $T'_0 > 0$  on  $[-1/2, 1/2] \setminus \{0\}$ ;
- 4  $\pm 1/2$  are preperiodic repelling for  $T_0$ .

# Family of 2-almost persistent attractors

Rovella (Bull. Braz. Math. Soc., 1993) showed that the flow of this vector field  $G_0$  has an attractor

$\Lambda_0 = \overline{\cap_{t>0} \phi_t^{G_0}(U)}$  and **its perturbations admit a two-parameter family of vector fields which is “almost persistent”**, as follows.

There exists a 2-dimensional  $C^\infty$  submanifold  $N$  of  $C^3$  vector fields  $\mathfrak{X}^3(\mathbb{R}^3)$  containing  $G_0$  s.t. the subset  $S \subset N$  corresponding to an attractor  $\Lambda_{G_s} = \overline{\cap_{t>0} \phi_t^{G_s}(U)}$  for each  $s \in S$ , then

$$\lim_{r \rightarrow 0} \frac{\text{Leb}(B_r(x) \cap S)}{\text{Leb}(B_r(x))} = 1,$$

where  $B_r(x)$  is an  $r$ -ball in  $N$  and  $\text{Leb}$  is the area measure.

# Persistent asymptotic sectional-hyperbolicity

Theorem (Vivas, San Martin: Nonlinearity, 2020)

The attractor  $\Lambda_0$  is 2-dimensionally almost persistent **asymptotically sectional hyperbolic** in the  $C^3$  topology.

A compact invariant partially hyperbolic set  $\Lambda$  of a vector field  $G$ , with  $d_{cu} = 2 = \dim E^{cu}$ , whose singularities are hyperbolic, is *asymptotically sectional hyperbolic* if there exists  $c_* > 0$  so that

$$\limsup_{T \rightarrow \infty} \frac{1}{T} \log |\det(D\phi_T | E_x^{cu})| \geq c_*, \quad \forall x \in \Lambda \setminus \bigcup_{\sigma \in \Lambda \cap \text{Sing}(G)} W^s(\sigma).$$

# Statistical stability of the Rovella family

## Theorem (A., JSP, 2021)

The family  $\mathcal{G}$  of contracting Lorenz attractors, with trapping region  $U$ , is such that each of its elements admits a unique physical measure, whose basin covers  $U$  except for zero Leb-measure subset and is statistically stable.

The existence of the unique physical/SRB measure  $\mu_a$  for each  $G_a \in \mathcal{G}$  follows from the fact that

- the Poincaré map  $R_a(x, y) = (T_a(x), H_a(x, y))$  satisfies
  - $H_a(x, \cdot)$  is a uniform contraction;
  - $T_a$  is a one-dimensional non-unif. exp. map with slow recurrence to the discontin. critical point  $\{0\}$

then every ergodic acip  $\nu_a$  w.r.t.  $T_0$  induces  $\mu_a$  which is an ergodic hyperbolic SRB-measure w.r.t.  $G_a$  on  $\Lambda_a$ .

# The physical/SRB measure in the Rovella family

This ensures, by well-known arguments, that  $\mu_a$  **admits an absolutely continuous disintegration along unstable manifolds** and is an **ergodic physical measure**.

Moreover, since the flow direction on partially hyperbolic sets is contained in the central-unstable direction, then Oseledec's Theorem ensures

$$\int \log |\det(D\phi_1^{G_a} | E^{cu})| d\mu_a = \lambda^+(x) \geq c_* > 0,$$

where  $\lambda^+(x) = \lim_{T \rightarrow \infty} \log |\det(D\phi_T | E_x^{cu})|^{1/T}$  is the largest Lyap. exponent along the two-dimensional bundle  $E^{cu}$  for  $\mu_a$ -a.e.  $x$ . Hence (Ledrappier-Young characterization)

$$h_{\mu_a}(\phi_1^{G_a}) = \int \log |\det(D\phi_1^{G_a} | E^{cu})| d\mu_a > 0.$$

# Robust expansiveness

Denote by  $S(\mathbb{R})$  the set of surjective increasing continuous functions  $h : \mathbb{R} \rightarrow \mathbb{R}$ . The flow is **expansive on an invariant compact set**  $\Lambda$  if for every  $\varepsilon > 0$  there is  $\delta > 0$  s.t. for any  $h \in S(\mathbb{R})$  and  $x \in \Lambda$

$$d(\phi_t(x), \phi_{h(t)}(y)) \leq \delta, \forall t \in \mathbb{R} \implies \\ \implies \exists t_0 \in \mathbb{R} \text{ s.t. } \phi_{h(t_0)}(y) \in \phi_{[t_0-\varepsilon, t_0+\varepsilon]}(x).$$

$\mathcal{G}$  is **robustly expansive** on  $\Lambda_s = \overline{\cap_{t>0} \phi_t^{G_s}(U)}$ ,  $s \in N \cap S$ , if  $\exists$  nbhd.  $V$  of  $s$  in  $N$  s.t.  $\forall \varepsilon > 0 \exists \delta > 0$  s.t. for  $u \in V \cap S$ ,  $x \in \Lambda_u = \overline{\cap_{t>0} \phi_t^{G_u}(U)}$  and  $h \in S(\mathbb{R})$ , then

$$d(\phi_t^{G_u}(x), \phi_{h(t)}^{G_u}(y)) \leq \delta, \forall t \in \mathbb{R} \implies \\ \implies \exists t_0 \in \mathbb{R} \text{ s.t. } \phi_{h(t_0)}^{G_u}(y) \in \phi_{[t_0-\varepsilon, t_0+\varepsilon]}^{G_u}(x).$$

# Robust entropy expansiveness from robust expansiveness

## Proposition (Bowen, 1972)

A robustly expansive attracting set  $\Lambda_G(U)$  on a family  $\mathcal{G} : N \rightarrow \mathcal{X}^r(M)$  admits  $\delta > 0$  which is a constant of  $h$ -expansiveness for each flow in the family.

**To fulfill all the conditions of statistical stability, it is enough to obtain**

## Lemma (Robust expansiveness for Rovella attractors)

The family  $\mathcal{G}$  of Rovella attractors is robustly expansive.

# Robust expansiveness for Rovella attractors

This is a consequence of the **locally eventually onto** property as follows. We write  $c_a^\pm = T_a(0^\pm) = \lim_{t \rightarrow 0^\pm} f(t)$  and note that  $c_a^- < 0 < c_a^+$  and  $c_a^\pm \rightarrow \pm 1/2$  when  $a \rightarrow 0$ .

**Lemma l.e.o. (Lemma 4.1 in Metzger: An. l'IHP, 2000)**

There exists a  $C^3$  neighborhood  $\mathcal{V}$  of  $G_0$  so that if  $G_a \in \mathcal{V}$ , then the map  $T_a$  is locally eventually onto, that is, for any interval  $J \subset [-1/2, 1/2] \setminus \{0\}$  there exists  $n = n(J) > 0$  so that  $f^n(J) \subset [c_a^-, c_a^+]$ .

Consequently, **there does not exist a pair of points**  $x_0 < y_0$  **with the same sign in**  $[-1/2, 1/2] \setminus \{0\}$  **so that**  $T_a^n[x_0, y_0]$  **does not contain the origin for all**  $n \geq 1$ .



# Proof of statistical stability

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## Consequence of continuity of equilibrium states

# Argument for statistical stability

## Theorem (Continuity of equilibrium states)

Let  $f : X \times M \rightarrow M$  and  $\psi : X \times M \rightarrow \mathbb{R}$  be continuous, with  $X = N \cap S$ , s.t.

- 1  $f_a$  admits some equilibrium state for  $\psi_a$ :  
 $\exists \mu_a \in \mathcal{P}_{f_a}(M)$  s.t.  $P_{f_a}(\psi_a) = h_{\mu_a}(f_a) + \int \psi_a d\mu_a$  for all  $a \in X$ .
- 2 For each weak\* accumulation point  $\mu_0$  of  $\mu_a$  when  $a \rightarrow * \in X$ , let  $a_k \rightarrow *$  s.t.  $\mu_k = \mu_{a_k} \rightarrow \mu_0$ , write  $f_k = f_{a_k}$ ,  $\psi_k = \psi_{a_k}$  and assume also that
  - 1 there exists a finite Borel partition  $\xi$  of  $M$  such that  $h_{\mu_k}(f_k) = h_{\mu_k}(f_k, \xi)$  for all  $k \geq 1$ ; and  $\mu_0(\partial\xi) = 0$ .
  - 2  $P_{f_k}(\psi_k) \rightarrow P_{f_*}(\psi_*)$  when  $k \rightarrow \infty$ .

Then every weak\* accumulation point  $\mu$  of  $(\mu_k)_{k \geq 1}$  when  $k \rightarrow \infty$  is a equilibrium state for  $f_*$  and the potential  $\psi_*$ .

# Entropy expansiveness and generating partitions

## Theorem (Bowen, 1972)

Let  $M$  be a compact metric space of finite dimension,  $\varepsilon > 0$  an  $h$ -expansiveness constant for  $f : M \rightarrow M$ , and  $\xi$  a Borel partition of  $M$  with  $\text{diam}(\xi) < \varepsilon$ . Then  $h_\mu(f) = h_\mu(f, \xi)$  for each  $f$ -inv. prob. measure  $\mu$ .

We may now use the robust entropy expansiveness to build a uniform generating partition  $\xi$  satisfying the conditions for continuity of equilibrium states, together with the assumption

$$h_{\mu_a}(\phi_1^{G_a}) - \int \log |\det(D\phi_1^{G_a} | E^{cu})| d\mu_a = 0$$

to apply the Theorem on Cont. of Eq. States with  $\psi_a = -\log |\det(D\phi_1^{G_a} | E^{cu})|$ .

**THE END.**

**OBRIGADO!**

# PARABÉNS ZÉZE!

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